

Sterile neutrinos at the eV scale

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Hinchcliffe's theorem

“When a title is in the form of a question,
the answer is always NO.”

see, however:

IS HINCHLIFFE'S RULE TRUE? •

Boris Peon

Abstract

Hinchliffe has asserted that whenever the title of a paper is a question with a yes/no answer, the answer is always no. This paper demonstrates that Hinchliffe's assertion is false, but only if it is true.

Recap: summary of neutrino oscillation results



- Established theoretical formalism
- Precise measurements of Δm_{23}^2 , $j\Delta m_{31}^2j$, Δm_{12}^2 , Δm_{21}^2 .



Recap: summary of neutrino oscillation results



- $\text{sgn}(\Delta m_{31}^2)$ unknown
- No sensitivity to CP yet
- Absolute neutrino masses not known
- Some open questions regarding coherence properties of neutrinos



Recap: summary of neutrino oscillation results



- LSND and MiniBoonE
 - ▶ Anomalous $\stackrel{\leftrightarrow}{e}$ appearance at short baseline
 - Reactor and gallium anomalies
 - ▶ Anomalous $\stackrel{\leftrightarrow}{e}$ disappearance at short baseline
- / Today's lecture



Oscillation anomalies: LSND and MiniBooNE

- LSND:

- ▶ e appearance in μ beam from stopped pion source (> 3) at $L=E$ 1 km/GeV

- MiniBooNE:

- ▶ No significant e or e excess in the LSND-preferred region
- ▶ but e consistent with LSND
- ▶ Low- E excess not understood

ν ! ν_e oscillations at $L/E = 1 \text{ km/GeV}$?

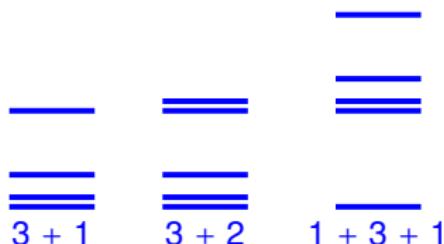
- Remember: Oscillation maxima for standard oscillations expected at
 - ▶ $L=E = 500 \text{ km/GeV}$ (from $m_{31}^2 = 2.4 \cdot 10^{-3} \text{ eV}^2$)
 - ▶ $L=E = 15\,000 \text{ km/GeV}$ (from $m_{21}^2 = 8.1 \cdot 10^{-5} \text{ eV}^2$)
- Explaining LSND and MiniBooNE requires an additional mass squared difference $\Delta m_{41}^2 = 1 \text{ eV}^2$.
- This requires an additional neutrino species.
- LEP measurements of the invisible Z width constrain the number of active neutrinos to three.
- Only possibility: A sterile neutrino ν_s , not coupling to SM gauge interactions.
 - ▶ “3 + 1 scenario”
- Then: Possibility of $\mu ! \nu_s ! \nu_e$ oscillations at $L=E = 1 \text{ km/GeV}$

ν_e appearance in the 3+1 scenario and beyond



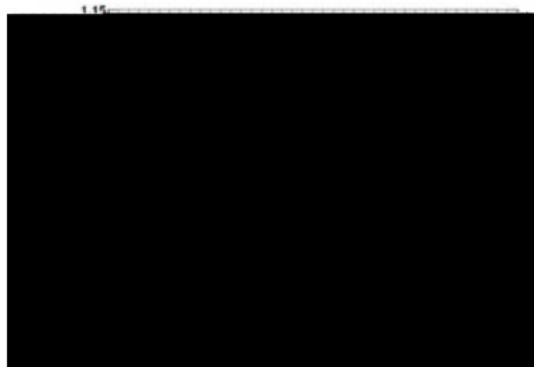
| | $\frac{2}{3+1}/\text{dof}$ | $\frac{2}{3+2}/\text{dof}$ | $\frac{2}{1+3+1}/\text{dof}$ |
|----------|----------------------------|----------------------------|------------------------------|
| LSND | 11.0/11 | 8.6/11 | 7.5/11 |
| MiniB | 19.3/11 | 10.6/11 | 9.1/11 |
| MiniB | 10.7/11 | 9.6/11 | 12.7/11 |
| E776 | 32.4/24 | 29.2/24 | 31.3/24 |
| KARMEN | 9.8/9 | 8.6/9 | 9.0/9 |
| NOMAD | 0.0/1 | 0.0/1 | 0.0/1 |
| ICARUS | 2.0/1 | 2.3/1 | 1.5/1 |
| Combined | 87.9/66 | 72.7/63 | 74.6/63 |

- Global fit to all appearance data is consistent
- Background oscillations important in MiniBooNE and E776
- Significant improvement in $3+2$ and $1+3+1$



The Gallium anomaly

- Intense radioactive e^- sources (^{51}Cr and ^{37}Ar) have been deployed in the GALLEX and SAGE solar neutrino detectors
- Neutrino detection via $^{71}\text{Ga} + e^- \rightarrow ^{71}\text{Ge} + e^-$
- Result: Measurements consistently lower than expectation (2.7%)



Giunti Laveder arXiv:1005.4599, arXiv:1006.3244
Mention et al. Moriond 2011 talk

- Question: How well are efficiencies of the radiochemical method understood?

The reactor anomaly

- Recent reevaluation of expected reactor $\bar{\nu}_e$ flux is **3.5% higher** than previous prediction Mueller et al. arXiv:1101.2663, confirmed by P. Huber arXiv:1106.0687
- Method:** Use measured $\bar{\nu}$ -spectra from ^{238}U , ^{235}U , ^{241}Pu fission at ILL and convert to $\bar{\nu}_e$ spectrum (for single $\bar{\nu}$ -decay: $E_\nu = Q - E_e$)
- Problem:** Requires knowledge of Q -values for **all** contributing decays.
! take from nuclear databases where available, fit to data otherwise

| Old method | New method |
|--|---|
| Schreckenbach 1985 30 effective decays (fit parameters to ILL data) | Mueller et al. arXiv:1101.2663 Uses nuclear databases (90% of $\bar{\nu}_e$ flux) 5 effective decays (remaining 10%) Error propagation, correlation matrix Corrections to the Fermi theory of decay Off-equilibrium corrections (not all $\bar{\nu}$ -decay chains in equilibrium) |

- Cross check:**
 - Simulate **mock e^- spectra** using few well-understood $\bar{\nu}$ -decays
 - Reconstruct $\bar{\nu}_e$ spectrum using **old method**: Result is **3% too low**
 - Reconstruct $\bar{\nu}_e$ spectrum using **new method**: Result is **exact**.

The reactor anomaly

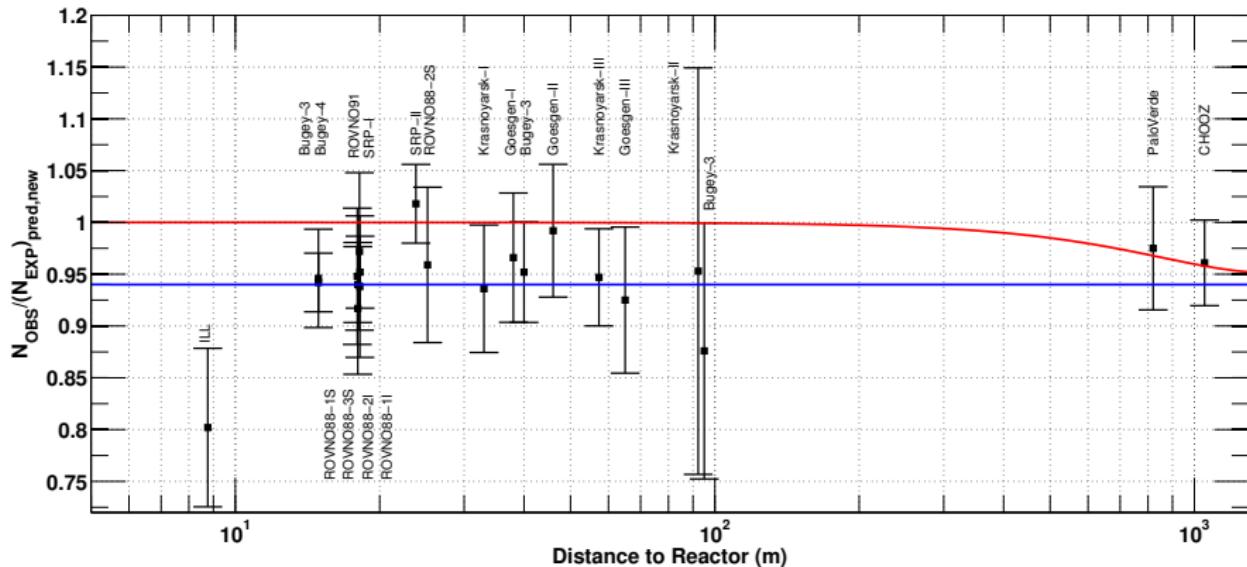
- Recent reevaluation of expected reactor $\bar{\nu}_e$ flux is **3.5% higher** than previous prediction Mueller et al. arXiv:1101.2663, confirmed by P. Huber arXiv:1106.0687
- Method:** Use measured γ -spectra from ^{238}U , ^{235}U , ^{241}Pu fission at ILL and convert to $\bar{\nu}_e$ spectrum (for single β -decay: $E_\nu = Q - E_e$)
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- Possible problems:**
 - Poorly understood effects in nuclei with large $\log ft$ Huber arXiv:1106.0687
 - Large systematic uncertainties for **non-unique forbidden decays** Hayes et al. arXiv:1309.4146

The reactor anti-neutrino anomaly

- Have short-baseline reactor experiments observed a $\bar{\nu}_e$ deficit?

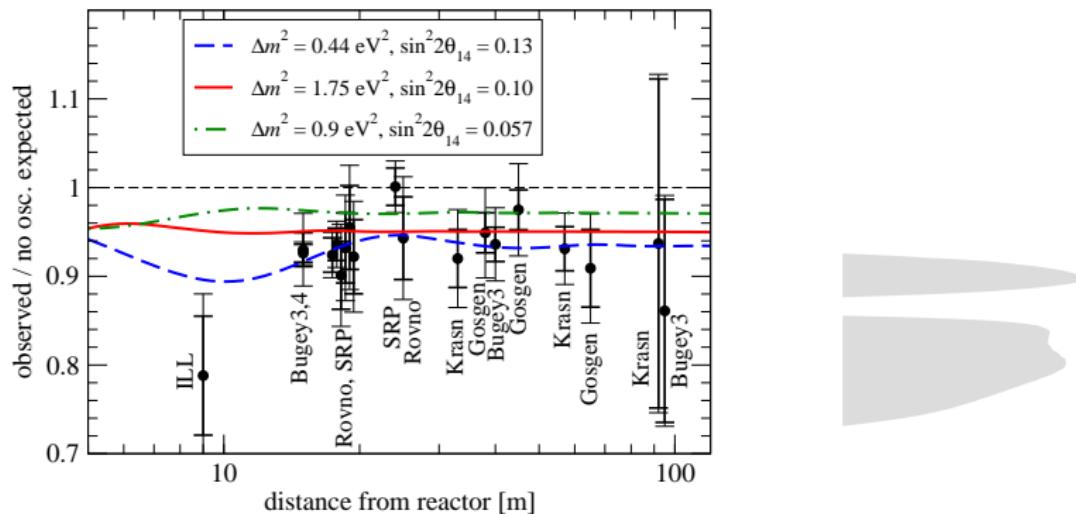


Mention et al. arXiv:1101.2755

red = new reactor $\bar{\nu}_e$ flux prediction

blue = old reactor $\bar{\nu}_e$ flux prediction

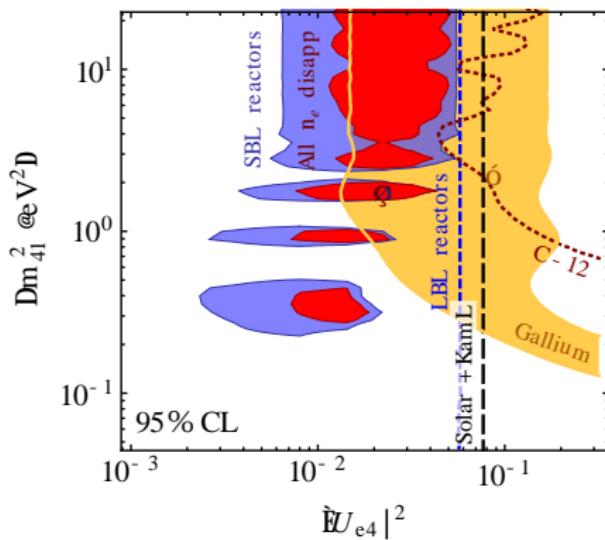
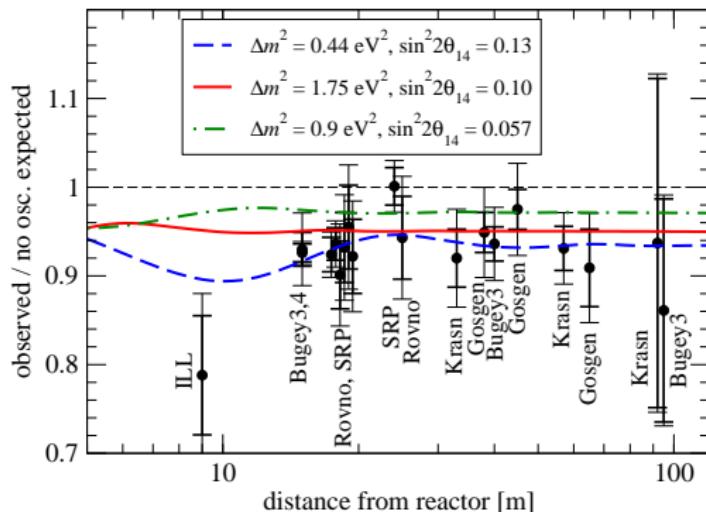
ν_e disappearance in the 3+1 scenario



| | $\sin^2 2\theta_{14}$ | $m_{41}^2 [\text{eV}^2]$ | $\chi^2_{\text{min}}/\text{dof} (\text{GOF})$ | $\chi^2_{\text{no osc}}/\text{dof} (\text{CL})$ |
|-------------------------|-----------------------|--------------------------|---|---|
| SBL rates only | 0.13 | 0.44 | 11.5/17 (83%) | 11.4/2 (99.7%) |
| SBL incl. Bugey3 spect. | 0.10 | 1.75 | 58.3/74 (91%) | 9.0/2 (98.9%) |

JK Machado Maltoni Schwetz, 1303.3011

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| SBL + Gallium | 0.11 | 1.80 | 64.0/78 (87%) | 14.0/2 (99.9%) |
| global $\bar{\nu}_e$ disapp. | 0.09 | 1.78 | 403.3/427 (79%) | 12.6/2 (99.8%) |

JK Machado Maltoni Schwetz, 1303.3011

Relation between appearance and disappearance

We find: $\stackrel{\leftarrow}{e}$ disappearance experiments consistent among themselves, $\stackrel{\leftarrow}{e}$ appearance experiments consistent among themselves.

But:

3 + 1 neutrinos

At large baseline ($L \gg 4E/\Delta m_{41}^2$), but $L \ll 4E/\Delta m_{31}^2$

$$P_{ee} = 1 - 2|U_{e4}|^2(1 - |U_{e4}|^2)$$

$$P_{\mu\mu} = 1 - 2|U_{\mu 4}|^2(1 - |U_{\mu 4}|^2)$$

$$P_{e\mu} = 2|U_{e4}|^2|U_{\mu 4}|^2$$

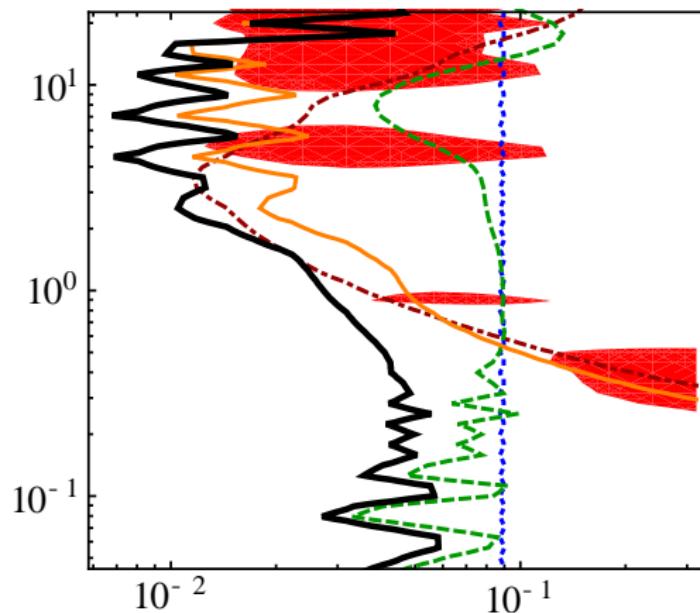
It follows

$$2P_{e\mu} = (1 - P_{ee})(1 - P_{\mu\mu})$$

In the **3 + 1** case, at **large enough baseline**, there is a **one-to-one relation** between the **appearance and disappearance probabilities**.

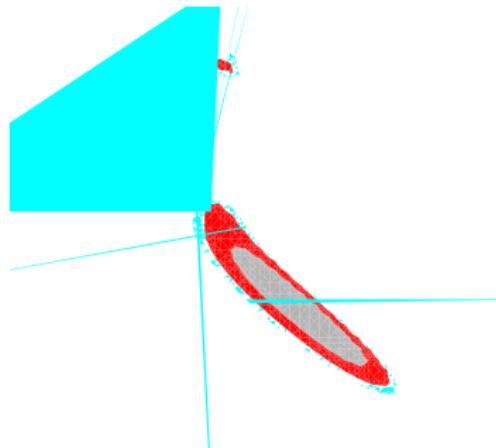
ν disappearance in the 3+1 scenario

- Parameter regions favored by **tentative hints** are in tension with null results from μ disappearance searches



The global oscillation fit

JK Machado Maltoni Schwetz, 1303.3011



| | $\chi^2_{\text{min}}/\text{dof}$ | GOF |
|-----|----------------------------------|-----|
| 3+1 | 712/(689 - 9) | 19% |

The global oscillation fit

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3 + 1 Severe **tension** between
appearance and
disappearance and between
exp's with and without a signal

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Parameter goodness of fit (PG) test:

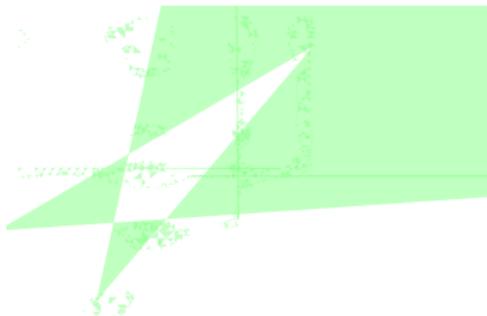
Compares $\frac{\chi^2_{\text{min}}}{\text{dof}}$ from global and separate
fits to test **compatibility** of 2 data sets

| | $\frac{\chi^2_{\text{min}}}{\text{dof}}$ | GOF | $\frac{\chi^2_{\text{PG}}}{\text{dof}}$ | PG |
|-----|--|-----|---|---------------------|
| 3+1 | 712/(689 - 9) | 19% | 18.0/2 | 1.2 $\cdot 10^{-4}$ |

The global oscillation fit

JK Machado Maltoni Schwetz, 1303.3011

- 3 + 1 Severe **tension** between appearance and disappearance and between exp's with and without a signal
- 3 + 2 Fit improves considerably with two sterile neutrinos



Parameter goodness of fit (PG) test:

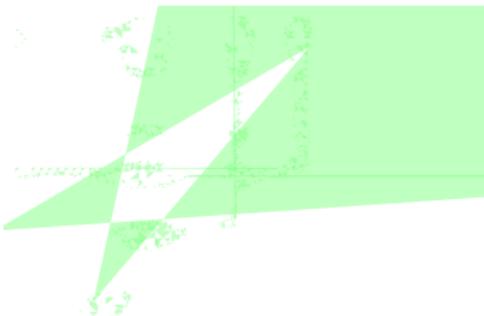
Compares χ^2_{min} from global and separate fits to test compatibility of 2 data sets

| | $\chi^2_{\text{min}}/\text{dof}$ | GOF | $\chi^2_{\text{PG}}/\text{dof}$ | PG |
|-----|----------------------------------|-----|---------------------------------|---------------|
| 3+1 | 712/(689 9) | 19% | 18.0/2 | 1.2 10^{-4} |
| 3+2 | 701/(689 14) | 23% | 25.8/4 | 3.4 10^{-5} |

The global oscillation fit

JK Machado Maltoni Schwetz, 1303.3011

- 3 + 1 Severe **tension** between appearance and disappearance and between exp's with and without a signal
- 3 + 2 Fit improves considerably with two sterile neutrinos
- 1 + 3 + 1 Further improvement, especially in **appearance** fit



Parameter goodness of fit (PG) test:

Compares χ^2_{min} from global and separate fits to test compatibility of 2 data sets

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| 3+1 | 712/(689 9) | 19% | 18.0/2 | 1.2 10^{-4} |
| 3+2 | 701/(689 14) | 23% | 25.8/4 | 3.4 10^{-5} |
| 1+3+1 | 694/(689 14) | 30% | 16.8/4 | 2.1 10^{-3} |

Conclusion from oscillation fits:

severe tension

in all cases

Sterile neutrinos in cosmology

Models with $\mathcal{O}(\text{eV})$ sterile neutrino(s) constrained by cosmology:

Sum of neutrino masses

$$\sum m_\nu \sim 0.23 \text{ eV}$$

of relativistic species

$$N_\nu = 4 \text{ mildly disfavored}$$

Ade et al. (Planck), arXiv:1303.5076
Gonzalez-Garcia Maltoni Salvado, arXiv:1006.3795
Hamann Hannestad Raffelt Tamborra Wong, arXiv:1006:5276

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Question:

What does it take to **evade** these constraints?

Are light sterile neutrinos ruled out by cosmology?

s production in the early Universe through e, μ, τ ! s oscillations at T & MeV

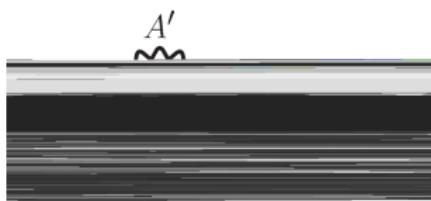
Dodelson Widrow 1994

Reconciling sterile neutrinos with cosmology

- Large lepton asymmetry (& 0.01) ! s production MSW-suppressed
Foot Volkas hep-ph/9508275, Chu Cirelli astro-ph/0608206, Saviano et al. arXiv:1302.1200
- Very low reheating temperature
Gelmini Palomares-Ruiz Pascoli, astro-ph/0403323
- Entropy production after neutrino decoupling (e.g. due to late decay of heavy sterile neutrinos or other particles) ! neutrinos diluted
Fuller Kishimoto Kusenko 1110.6479, Ho Scherrer 1212.1689
- Self-interacting s
! s production MSW suppressed
Hannestad Hansen Tzanavaris arXiv:1310.5926; Dasgupta JK arXiv:1310.6337

Suppressed ν_s production from thermal MSW effect

- ν_s production in the early Universe through e, μ, τ ! ν_s oscillations
Dodelson Widrow 1994
- Assume ν_s couple to & MeV gauge boson A'
- Neutrino self energy:



- Leads to thermal MSW potential $V \propto T^4$
- At high T : mixing strongly suppressed

$$\sin 2 \theta_{\text{eff}} = \frac{\sin 2 \theta}{\sqrt{\sin^2 2 \theta + (\cos 2 \theta - \frac{2EV}{\Delta m^2})^2}}$$

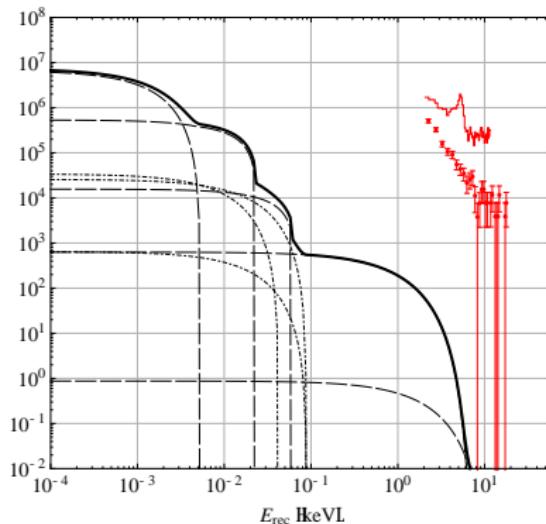
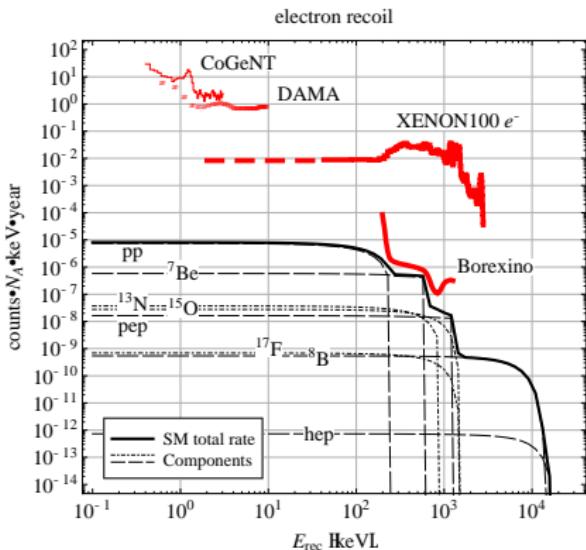
- No ν_s production through oscillations
! no cosmological constraints

Hannestad Hansen Tram arXiv:1310.5926
Dasgupta JK arXiv:1310.6337

Neutrinos and direct dark matter detection

Solar neutrinos are a well-known **background** to future direct DM searches:

see e.g. Gütlein et al. arXiv:1003.5530



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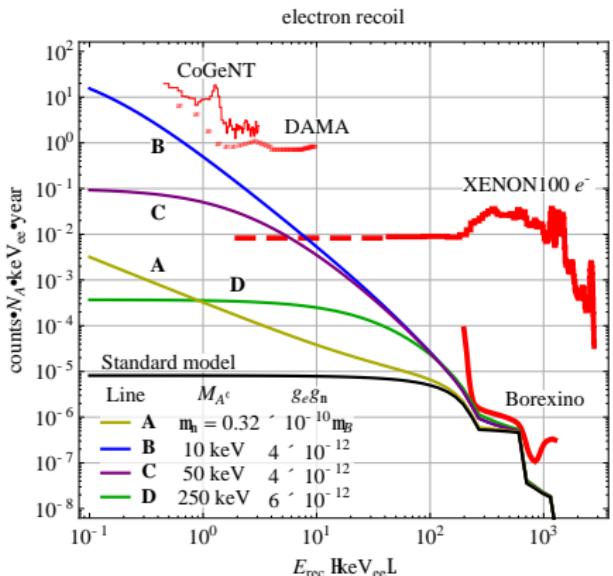
SM signal will only become sizeable in multi-ton detectors

But: New physics can enhance the rate

Examples:

- Neutrino magnetic moments
- Sterile neutrinos + . GeV scale hidden sector gauge force

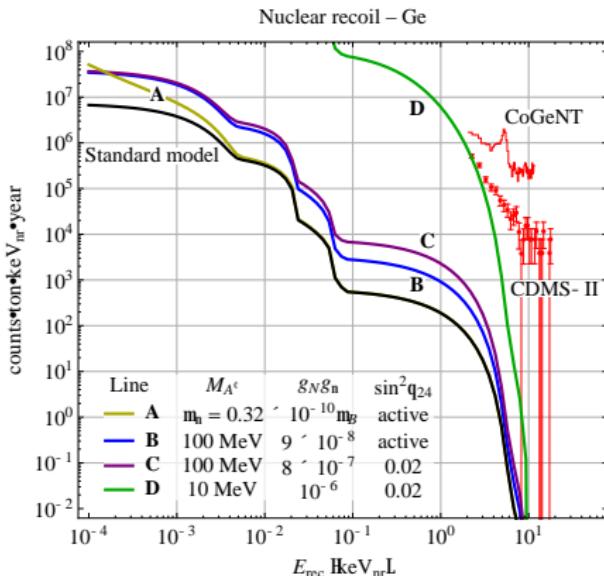
Low-energy scattering of neutrinos beyond the SM



A: magnetic moment

B, C, D: kinetically mixed A' + sterile s

- Enhanced scattering at low E_r for light A'
- Negligible compared to SM scattering ($\sim g^4 m_T = M_W^4$) at energies probed in dedicated neutrino experiments



A: magnetic moment

B: $U(1)_{B-L}$ boson

C: kinetically mixed $U(1)'$ + sterile

D: $U(1)_B$ + sterile charged under $U(1)_B$

[Pospelov 1103.3261, Pospelov Pradler 1203.0545]

[Harnik JK Machado arXiv:1202:6073]

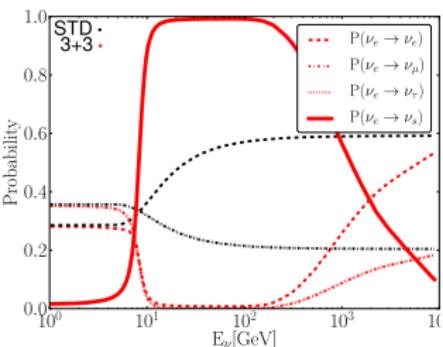
Sterile neutrinos and DM annihilation in the Sun

- Neutrino telescope limits on neutrinos from dark matter annihilation in the Sun depend crucially on oscillation physics.
- If sterile neutrinos exist, new MSW resonances can lead to strong conversion of active neutrinos into sterile neutrinos in the Sun.

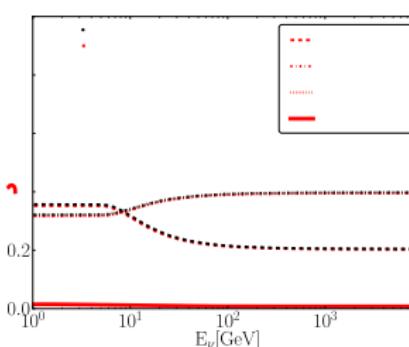
Esmaili Peres, arXiv:1202.2869
Argüelles JK, arXiv:1202.3431

Oscillation probabilities for a 3+3 toy scenario

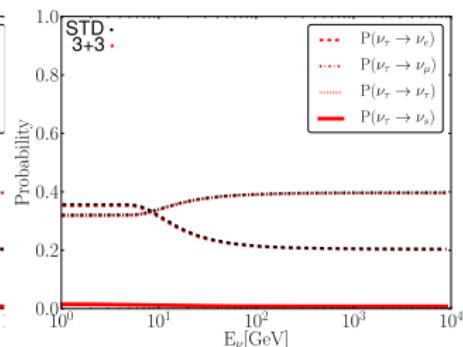
$P(e^- / x)$



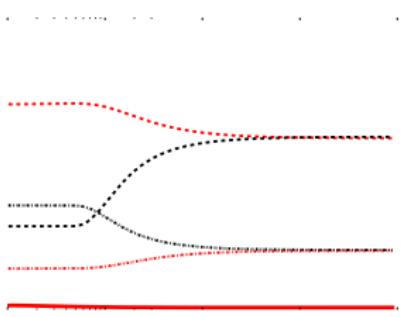
$P(\mu^- / x)$



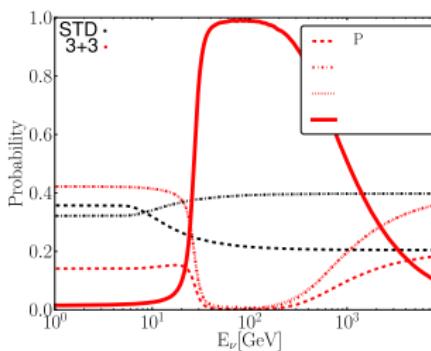
$P(\tau^- / x)$



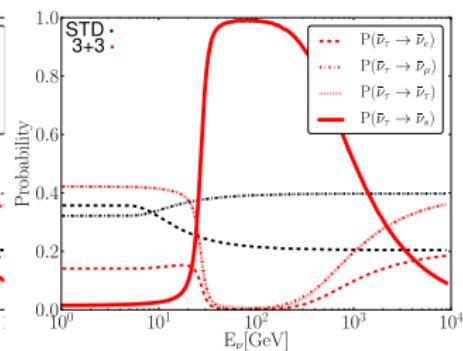
$P(\bar{e}^- / \bar{x})$



$P(\bar{\mu}^- / \bar{x})$



$P(\bar{\tau}^- / \bar{x})$



Thick red lines = active–sterile oscillations

plots from Argüelles JK, arXiv:1202.3431
see also Esmaili Peres, arXiv:1202.2869